Temperature Dependence of the Mössbauer Recoil-free Fraction as a Probe of the Lattice Structure of Tin Compounds

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Summary The temperature coefficient of the Mössbauer recoil-free fraction is a useful guide to the nature of the lattice structure of tin compounds.

DURING the past decade, Mössbauer spectroscopy has emerged as a particularly useful structural technique in tin chemistry. However, by far the majority of the accumulated data refer to the application of isomer shift and quadrupole splitting measurements to the solution of structure and bonding problems within the first co-ordination sphere of the metal atom. By measuring the variation with temperature of the recoil-free fraction, a measure of the number of nuclei undergoing recoil-less γ -ray resonance, Mössbauer spectroscopy may also be applied to the investigation of the lattice structure of tin-containing materials.

Since for a thin absorber the recoil-free fraction, f(T), is directly related to the area under the resonance curve, A(T), of a Mössbauer spectrum, the temperature variation of f(T) may be followed by evaluating the experimentally more accessible parameter, A(T). We have collected resonance area data for the compounds listed in the Table, and semilogarithmic plots of the resonance areas (normalised to 77 K) against temperature were found to be linear in the temperature range $77 \ll T \ll 150$ K, as expected TABLE. Recoil-free fraction temperature coefficients.

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|--|------------------|-----|----|------------|-------------------------------|---------------|--|
| Compound | | | | | $a	imes 10^2/\mathrm{K^{-1}}$ | Valance state | Comment on Lattice Structure |
| Me ₃ SnONPhCOPh | (1). | | | | 1.74. | IV | Crystals composed of monomeric molecules $(X-ray)^{\mathbf{s}}$ |
| Ph ₃ SnONPhCOPh | (2). | | | | 1·85 [°] | IV | b |
| Me ₃ SnO ₃ SPh·H ₂ O (| (3) . | • | •• | •• | 1.71 | IV | Adjacent molecules weakly linked by H-bonds to form a 1-D infinite chain $(X-ray)^{\circ}$ |
| $Me_{3}SnO_{2}CC_{5}H_{4}N-2$ | H ₂ O | (4) | •• | ••• | 1.27 | IV | Adjacent molecules linked by a 3-D network of H-bonds $(X-ray)^d$ |
| Me ₃ SnONC ₆ H ₁₀ (5) | | | | | 0.97 | IV | 1-D co-ordinate chain polymer (X-ray) ^e |
| Et ₃ SnONC ₆ H ₁₀ ⁽⁶⁾ | | | | | 1.16 | īv | 1 D to ordinate enant polymon (11 103) |
| Pr ⁿ ,SnONC,H ₁₀ (7) |) . | | | | 1.43 | ĪV | |
| PhaSnONC, H10 (8) | ′. | | | | 1.82. | ĪV | |
| Me ₂ Sn(ONMeCOMe | e), (9) | | | | 1.85° | ĪV | Crystals composed of monomeric molecules $(X-ray)^{\dagger}$ |
| Me ₂ Sn(ONHCOMe) | , (10) | | | | 0.92 | ĪV | Adjacent molecules linked by H-bonds (X-ray) ^g |
| ClMe,SnSnMe,Cl (1 | 1) . | | | | 1.93 | IV | |
| Me ₂ SnO (12) | <i>.</i> . | | | | 0.87 | IV | Polymeric |
| $[Sn^{IV}Sn^{1i}(O'_{a}CC_{e}H_{4}NO_{5}-o)_{4}O\cdot thf]_{s}^{-1}$ | | | | | | | |
| (a) Sn ^{IV} . | | | | <i>.</i> . | 1.30 | IV | Crystals composed of tetranuclear cluster molecules $(X-ray)^{h}$ |
| (b) Sn ^{II} | | | | | 1.70 | II | |
| $Sn(O_2CH)_2$ (14) | | | | • • | 0.84 | II | 2-D Sheet polymer (X-ray) |
| SnO (15) | | | | | 0.23 | II | Ionic lattice (X-rav) ^k |
| SnCl ₂ (16) | | | | | 0.92 | II | PbCl ₂ -type lattice $(X - ray)^1$ |
| $Sn(C_{5}H_{5})_{2}(17)$ | • | • | •• | •• | 3.13 | II | · 2 · 51 · · · · · · · · · · · · · · · · · |
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according to the Debye model of solids in the high temperature limit. The temperature coefficients of the recoil-free fraction (a/K^{-1}) derived therefrom are also listed in the Table, together with other available data concerning the solid-state structure for the materials.

From a comparison with those compounds of known structure, it is apparent that those organotin(IV) compounds which have lattices consisting of non-interacting monomeric molecules [(1), (2), and (9)] have temperature coefficients of ca. $1.8 \times 10^{-2} \text{ K}^{-1}$, irrespective of the coordination number at tin. Weak intermolecular interactions, such as the hydrogen-bonding between adjacent $Me_{3}SnO_{3}SPh \cdot H_{2}O$ molecules (3) reduces the value of a slightly (to $1.71 \times 10^{-2} \text{ K}^{-1}$), whilst a more complex system of hydrogen bonds as in $Me_3SnO_2CC_5H_4N-2\cdot H_2O$ (4) causes a larger reduction (to $1.27 \times 10^{-2} \,\mathrm{K^{-1}}$), reflecting the increasing rigidity of the lattice. Strongly hydrogenbonded lattices such as (10), and one-, two-, and threedimensional polymeric lattices [e.g. (5), (14), and (12),respectively] exhibit temperature coefficients of ca. 0.9 $\times 10^{-2}$ K⁻¹, as does tin(II) chloride (16), which possesses the PbCl₂-type lattice. Tin(II) oxide (15) exhibits the lowest temperature coefficient of all the compounds measured, 0.23×10^{-2} K⁻¹, consistent with its tightly packed ionic lattice.

Using these data, we can thus infer that $Ph_3SnONC_8H_{10}$ (8) $(a = 1.825 \times 10^{-2})$ is unassociated, and that the ethyl and n-propyl homologues [(6) and (7), respectively] with a values of 1.16×10^{-2} and 1.43×10^{-2} K⁻¹, respectively, have structures intermediate between the one-dimensional co-ordinate polymeric structure determined for the methyl homologue (5) and the monomeric phenyl derivative (8), demonstrating the weakening of the co-ordinate interaction as the steric bulk of the organic groups bonded to tin increases. The high value for 1,2-dichloro-1,1,2,2-tetramethylditin (11) $(1.93 \times 10^{-2} \text{ K}^{-1})$ is also indicative of the absence of significant intermolecular interaction for this compound in the solid. Dicyclopentadienyltin(II) (17) exhibits the highest temperature coefficient recorded $(3.13 \times 10^{-2} \text{ K}^{-1})$, being higher than that reported for bis(propane-1,3-dionato)tin(II) $(2.8 \times 10^{-2} \text{ K}^{-1})$,¹ which is most probably isostructural with bis(1-phenylpropane-1,3dionato)tin(II)² Such values of *a* are clearly consistent with the presence of non-interacting monomeric units in the solid for these materials [cf. the structure of dicyclopentadienyl-lead(II) in the solid³].

The values of a for the two types of tin atom in the mixed-valence carboxylate (13) are interesting. Crystals of the compound consist of tetranuclear clusters with the Sn^{IV} atoms tightly bound in an octahedral arrangement, and the Sn¹¹ atoms enjoying pseudo-pentagonal bipyramidal co-ordination.4 The temperature coefficient values of $1.30 \times 10^{-2} \, \mathrm{K^{-1}}$ for the SnIV atoms and $1.70 \times 10^{-2} \, \mathrm{K^{-1}}$ for the Sn^{II} atoms thus reflect the oligomeric constitution of the material, as well as the relative strengths of binding of the two tin atoms within the macromolecule (Sn $^{\rm IV}$ > Sn^{II})

We thank the S.R.C. for support.

(Received, 30th May 1977; Com. 532.)

 $\dagger a = - d [\ln A(T)/A(77)]/dT.$

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